# Hexagonal Tungsten Bronze-Type $Fe^{III}$ Fluoride: $(H_2O)_{0.33}FeF_3$ ; Crystal Structure, Magnetic Properties, Dehydration to a New Form of Iron Trifluoride

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(H<sub>2</sub>O)<sub>0.33</sub>FeF<sub>3</sub>, grown by hydrothermal synthesis, crystallizes in the orthorhombic system with cell dimensions a = 7.423(3) Å, b = 12.730(4) Å, c = 7.526(3) Å, and space group *Cmcm*, Z = 12. The structure, derived from single crystal X-ray diffraction data (605 independent reflections) is refined to R = 0.019 ( $R_{\omega} = 0.021$ ). The framework of the Fe<sup>III</sup>F<sub>6</sub> octahedra is related to that of hexagonal tungsten bronze (HTB) Rb<sub>0.29</sub>WO<sub>3</sub>. At 122°C, zeolithic water is evolved from hexagonal tunnels without any noticeable change of the fluorine skeleton. The related anhydrous compound represents a new form of iron trifluoride which is denoted HTB-FeF<sub>3</sub>; at 525°C, it transforms into the cubic form of ReO<sub>3</sub>-type. (H<sub>2</sub>O)<sub>0.33</sub>FeF<sub>3</sub> and HTB-FeF<sub>3</sub> are antiferromagnetic, with Néel temperatures of  $T_N = 128°7 \pm 0.5$  K and  $T_N = 97 \pm 2$  K, respectively.

#### Introduction

During the thermal decomposition of tetragonal  $\beta$ -FeF<sub>3</sub>, 3 H<sub>2</sub>O, Macheteau and Charpin (1) isolated an intermediate hydrate to which they assigned the formula FeF<sub>3</sub>, H<sub>2</sub>O. Its X-ray spectrum was indexed in a hexagonal cell with parameters a= 7.395 Å, c = 7.557 Å; however, the product always contained some amounts of  $\alpha$ -Fe<sub>2</sub>O<sub>3</sub> and  $\alpha$ -FeF<sub>3</sub> · 3 H<sub>2</sub>O. This fact cast some doubt on the chemical composition given by these authors and led us to reexamine this compound.

Hydrothermal conditions were systematically applied to  $\text{FeF}_3(2)$  in order to render the previous results more precise: anhydrous or hydrated phases appear, depending on the temperature and fluorine concentration of the solvent; the trihydrate  $\beta$ -FeF<sub>3</sub> · 3 H<sub>2</sub>O is obtained at 250°C from aqueous solutions, while at 380°C in 49% HF, anhydrous FeF<sub>3</sub> crystals are grown. For intermediate conditions, a new compound (H<sub>2</sub>O)<sub>0.33</sub>FeF<sub>3</sub> is isolated. Its structure is related to that of the hexagonal tungsten bronze (HTB) Rb<sub>0.29</sub>WO<sub>3</sub> (3).

### **Preparation and Characterization**

Light-green crystals of  $(H_2O)_{0.33}$ FeF<sub>3</sub> are grown by hydrothermal synthesis (FeF<sub>3</sub>: 1.2 g, H<sub>2</sub>O: 2.2 cm<sup>3</sup>) at 360°C, 200 MPa, for 4 days. They present a pseudo-hexagonal prismatic habit and frequent (110) twinning. Their powder X-Ray diffraction pattern is indexed in an ideal hexagonal bronze-type cell a = 7.39 Å, c = 7.56 Å; the orthorhombic deformation, described later, is not observed.

ATG experiments show the easy dehydration of (H<sub>2</sub>O)<sub>0.33</sub>FeF<sub>3</sub> at 122°C (heating rate:  $2^{\circ}$ C min<sup>-1</sup>). The weight loss, 5.3%, is close to theorical loss, 5.1%, corresponding to 0.33  $H_2O$  per 1 FeF<sub>3</sub> molecule. The resulting anhydrous compound represents a new form of iron trifluoride and will be denoted HTB-FeF<sub>3</sub>. X-Ray diffraction spectra of  $HTB-FeF_3$ , from room temperature to 300°C, are very similar to those of  $(H_2O)_{0.33}$ FeF<sub>3</sub>: only very small intensity differences are observed; lattice spacings are identical. At higher temperature, hydrolysis begins to occur on the heating stage. However, thermal analysis (DTA, DSC) indicates that at 60°C, HTB-FeF<sub>3</sub> undergoes an endothermic reversible phase transition, probably due to small reorientations of the iron octahedra. Furthermore, at 525°C under argon, an endothermic peak is observed: HTB-FeF<sub>3</sub> irreversibly transforms to ReO<sub>3</sub>-type FeFe<sub>3</sub>, cubic at this temperature, rhombohedral at  $T < 394^{\circ}C$  (4).

The infrared spectrum of  $(H_2O)_{0.33}FeF_3$ crystals shows the presence of intense and narrow peaks at 1625, 3550, and 3625 cm<sup>-1</sup> associated with weak and broad bands around 1600 and 3250 cm<sup>-1</sup>. The absorption close to 1600 and 3500 cm<sup>-1</sup> may be assigned, respectively, to the bending  $\delta_{H_2O}$ and to the symmetrical or antisymmetrical stretching  $\nu_{OH}$  vibrations. The presence at 1625 cm<sup>-1</sup> of an intense and narrow peak proves unquestionably that structural water is present as H<sub>2</sub>O molecules. If structural water were coordinated to iron atoms as hydroxyl groups, only  $\nu_{OH}$  would be expected. Due to the method of preparation, some adsorbed water, which contributes to the broad bands, is in evidence.

#### Experimental

Laue and precession photographs show that  $(H_2O)_{0.33}$ FeF<sub>3</sub> crystals are orthorhombic with space group Cmcm (No. 63) (Table I). The cell derives from the hexagonal bronze-type cell (dashed unit cell in Fig. 1: a = b = 7.370 Å,  $\gamma = 119^{\circ}52$ ). Intensity data are collected on a CAD 4 Nonius diffractometer.<sup>1</sup> Operating features are as follows:  $MoK\alpha$  radiation; graphite monochromator; scan-mode  $\omega$ ; sweep  $s = (0.80 + 0.35 \text{ tg } \theta)^{\circ}$ ; scintillation counter aperture D = (1.50 + 1.50)1.00 tan  $\theta$ ) mm; scanning speed v = $(20.1166/NV)^{\circ}$  min<sup>-1</sup> with NV integer; range measured:  $3^{\circ} \le \theta \le 40^{\circ}, 0 \le h \le + 14, -22$  $\leq k \leq +22, -12 \leq l \leq +12$  with C restrictions; reflections measured: total = 1778, mmm symmetry independent and  $|F| \ge 6\sigma$ (|F|) = 605. Corrections for Lorentz and polarization effects and for absorption are applied, using, for all calculations, the SHELX program (5). The transmission factor varies from 0.615 to 0.724. Atomic scat-

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 TABLE I

 Experimental Data for (H2O)0.33FeF3 Crystals

Symmetry: orthorhombic $Z = 12$	Cell parameters: $a = 7.423(3)$ Å
Space group: Cmcm	b = 12.730(4)  Å
Conditions: $hkl h + k = 2n$	c = 7.526(3)  Å
$h0l\ l=2n(h=2n)$	
$\rho_{exp} = 3.25(10) \text{ g cm}^{-3}$	$\mu(MoK\alpha) = 61.7 \text{ cm}^{-1}$
$\rho_{\rm calc} = 3.33 \ {\rm g \ cm^{-3}}$	Crystal dimensions: $(0.07 \times 0.10 \times 0.06)$ mm <sup>3</sup>

Z = 0.5



FIG. 1. (001) projection of  $(H_2O)_{0.33}FeF_3$ .

tering factors are taken from the "International Tables for X-Ray Crystallography" (1968) for  $Fe^{3+}$ , O, and  $F^{-}$  (6).

#### Structure Determination

Examination of Fe-Fe interactions in the Patterson map built from  $I_{hkl}$  clearly points out that iron atoms occupy a pseudo-hexagonal bronze-type array. Fe-F interactions further indicate that FeF<sub>6</sub> octahedra are tilted from the c axis; in particular, some of these octahedra strictly rotate around a. Thus, a solution is based on a model, giving to the  $Fe(1)F_6$  octahedron a tilt around a and correlated tilts to  $Fe(2)F_6$  octahedra. The water molecules are placed at the center of the hexagonal cavities. Refinement of this set of atomic coordinates, weighting, secondary extinction, and isotropic thermal parameters easily converges to R = 0.035 $(R_{\omega} = 0.041)$ . The residual falls to R = 0.019 $(R_{\omega} = 0.021)$  when applying anisotropic thermal motion for all nonhydrogen atoms.<sup>2</sup>

Table II presents the final results for the

<sup>2</sup> See NAPS document No. 04038 for 6 pages of supplementary material. Order from ASIS/NAPS, Microfiche Publications, P.O. Box 3513, Grand Central Station, New York, NY 10163. Remit in advance \$4.00 for microfiche copy or for photocopy, \$7.75 up to 20 pages plus \$.30 for each additional page. All orders seven independent positions; characteristic distances are listed in Table III. The refinement of the site occupancy factor for oxygen atoms in  $4a (0 \ 0 \ 0)$  or  $4c (0, y, \frac{1}{4})$  really indicates that the 4c site is fully occupied and the 4a site, empty.

#### **Discussion of the Structure**

The structure determination and the light-green color of the crystals prove that  $(H_2O)_{0.33}$ FeF<sub>3</sub> represents the first hexagonal tungsten bronze-type fluoride with only one metallic species and no mixed-metal valencies. Figure 1 presents a projection of the structure at z = 0.5 and shows its similarity to  $Cs_{0,4}$  Zn<sub>0,4</sub> Fe<sub>1,6</sub> F<sub>6</sub> (7). Iron atoms occupy two types of crystallographic sites. The six nearest-neighbor distances around iron atoms (Table III) are consistent with  $Fe^{3+}-F^-$  distances ( $\overline{d} = 1.947$  Å) and confirm the absence of H<sub>2</sub>O or OH groups in the iron coordination polyhedra (8). Refinement and Fourier map allow the placing of the oxygen atom of water molecules in the

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Atom	Site	X	y	2	$U_{11}\cdot 10^4$	$U_{22}\cdot 10^4$	$U_{33}\cdot 10^4$	$U_{23}\cdot 10^4$	$U_{13}\cdot 10^4$	$U_{12} \cdot 10^4$	Beq. (Å <sup>2</sup>
Fel	44	0	F	0	84(1)	65(1)	66(1)	-1(2)	0	0	0.57(1)
Fe2	8d	-14	-44	0	61(1)	88(1)	73(1)	-1(1)	2(1)	10(1)	0.58(2)
FI	8f	. 0	0.2156(1)	0.5511(2)	72(5)	244(8)	176(6)	17(6)	0	0	1.29(5)
F2	164	0.1876(1)	0.3962(1)	0.0434(2)	193(4)	124(4)	199(4)	-9(4)	-25(4)	62(4)	1.36(3)
- H	40	0	0.5386(2)	44	230(10)	133(10)	66(8)	0	0	0	1.13(7)
F4	200	0.2101(2)	0.2176(1)		179(7)	201(8)	78(6)	0	0	-2(7)	1.21(6)
H <sub>2</sub> O	4c	0	-0.0006(4)		519(28)	410(27)	721(40)	0	0	0	4.34(25)

TABLE II

TABLE III Characteristic Interatomic Distances (Å) and Bond Angles

$4 \times \text{Fe1}-\text{F2}$	1.948(1)	$\bar{d}_{\text{Fel}-\text{F}} = 1.947$	
$2 \times \text{Fe1}\text{-}\text{F3}$	1.945(1)	$d_{\rm F-F} = 2.753$	
$2 \times \text{Fe2}-\text{F1}$	1.945(1)	$\bar{d}_{\text{Fe2-F}} = 1.946$	
$2 \times \text{Fe2}-\text{F2}$	1.945(1)	$\hat{d}_{\rm E,E} = 2.752$	
$2 \times \text{Fe2}-\text{F4}$	1.949(1)	WF-F 2.752	
		Fe1-F3-Fe1	150°7(1)
$2 \times H_2O-F1$	3.119(4)	Fe2-F4-Fe2	149°8(1)
$2 \times H_2O-F1$	3.565(4)	Fe2-F1-Fe2	145°1(1)
$4 \times H_2O-F2$	3.086(5)	Fe1-F2-Fe2	142°3(1)
$4 \times H_2O-F2$	3.467(5)		
$2 \times H_2O-F3$	3.745(4)	F2-Fe1-F2	88°7(1)
$2 \times H_2O-F4$	3.186(5)	F2-Fe1-F3	89°5(1)
$2 \times H_{2}O-F4$	4.183(4)	F1-Fe2-F2	88°8(1)
2		F1-Fe2-F4	<b>89°9</b> (1)

hexagonal cavities, but hydrogen atoms cannot be located.

Orthorhombic symmetry arises mainly from anionic displacements. The nearly regular Fe(1)F<sub>6</sub> and Fe(2)F<sub>6</sub> octahedra tilt cooperatively from c axis (14°1 and 14°9, respectively). These important tilts imply a shortening of  $8H_2O-F$  distances (Table III). However, the shortest  $H_2O-F$  distances (3.09 Å) provide evidence that water molecules are weakly hydrogen bonded. They can move inside the large tunnels, as confirmed by the high values of their thermal motion parameters and by the easy dehydration of the compound.

Recently, Gerand *et al.* (9) synthetized  $WO_3 \cdot \frac{1}{3}H_2O$ . The structure, refined from powder diffraction data, is built up from bronze-type layers in a cell similar to that of  $(H_2O)_{0.33}FeF_3$ . However, the lattice-centering mode, F for  $WO_3\frac{1}{3}H_2O$  and C for  $(H_2O)_{0.33}FeF_3$ , implies that the connection of the layers is different for both compounds. In  $WO_3$ ,  $\frac{1}{3}H_2O$ , the layers are shifted along **a.** The water molecule, which belongs to the tungsten octahedron, points toward crossed tunnels running along [101] and [101]. In  $(H_2O)_{0.33}FeF_3$ , the layers are superimposed along **c.** 

Dehydration of  $(H_2O)_{0.33}$ FeF<sub>3</sub> and WO<sub>3</sub> ·  ${}^{1}_{3}$ H<sub>2</sub>O leads to similar HTB-type compounds. So their dehydration mechanisms must differ. For  $(H_2O)_{0.33}$ FeF<sub>3</sub>, the loss of zeolithic water is achieved at low temperature, and the framework of octahedra remains unchanged. By contrast, for WO<sub>3</sub> ·  ${}^{1}_{3}$ H<sub>2</sub>O, a shift of the HTB layers must occur; then, the transition is destructive and realized at high temperatures ( $T > 300^{\circ}$ C) (10).

### **Magnetic and Mossbauer Studies**

For (H<sub>2</sub>O)<sub>0.33</sub>FeF<sub>3</sub> crystals and HTB-FeF<sub>3</sub>, the thermal variations of inverse susceptibility  $\chi^{-1}$  (Fig. 2) and magnetization, measured with a vibrating sample magnetometer, are indicative of strong antiferromagnetic interactions. Even at 293 K, the Curie–Weiss law is not obeyed; so  $\theta_p$  values and molar Curie constants can not be obtained. The very flat  $\chi^{-1}$  versus T curves present a slight slope discontinuity around T = 130 K for  $(H_2O)_{0.33}$ FeF<sub>3</sub> crystals, T =100 K for anhydrous HTB-FeF<sub>3</sub>. Mössbauer experiments show that these temperatures correspond to the magnetic ordering temperatures  $T_N$  (128.7 ± 0.5 K for  $(H_2O)_{0.33}$ FeF<sub>3</sub> and 97  $\pm$  2 K for HTB-FeF<sub>1</sub>).

Experimental Mössbauer data are listed in Table IV. For  $(H_2O)_{0.33}$ FeF<sub>3</sub> crystals, the two types of iron sites are not distinguished



FIG. 2. Thermal variation of the magnetic inverse susceptibility  $(1/\text{emu g}^{-1})$ .

TABLE	IV
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Mössbauer Data

T(K)	δ <sup>a</sup> (mm s <sup>-1</sup> )	$\frac{\Delta E_{\rm Q}}{(\rm mm~s^{-1})}$	$\frac{4\epsilon^b}{(mm \ s^{-1})}$	H (kOe)	Г (mm s <sup>-1</sup> )
295	0.439(2)	0.640(4)	_	_	0.50(1)
129	0.539(2)	0.645(4)	—	_	0.52(1)
77	0.54(1)		-0.24(4)	511(2)	0.66(1)
4.2	0.55(1)		-0.26(4)	565(2)	0.60(1)

<sup>a</sup>  $\delta$  = isomer shift relative to Fe metal (T = 300 K).

<sup>b</sup>  $\varepsilon$  = quadrupolar shift of the outer Zeeman lines.

through Mössbauer spectrometry, as could be deduced from considerations of bond lengths (Table III) and angles. The thermal variation of their common hyperfine magnetic field is shown in Fig. 3; the value found at 4.2 K, close to that observed for hexagonal bronze-type  $A_x$ FeF<sub>3</sub> (11), corroborates the absence of water molecules in the iron coordination polyhedra. From the  $\varepsilon$  value, the determination of the angle between the magnetic hyperfine field direction and the main electric field gradient axis is ambigous: it could be 63°5 or 46°6, depending on the unknown sign of quadrupole splitting  $\Delta E_{\Omega}$ .

According to the Kanamori (12) or Goodenough rules (13), the antiferromagnetic exchange between layers is easily ra-

tionalized (Fe-F-Fe =  $149^{\circ}8$  or  $150^{\circ}7$ , Table III), but frustration must exist in the cycles formed by three octahedra linked by corners in the layer plane. Thus, neutron



FIG. 3. Thermal variation of the magnetic hyperfine field of  $(H_2O)_{0.33}FeF_3$ .

diffraction experiments, scheduled at the ILL Grenoble, will have to determine the magnetic structure of  $(H_2O)_{0.33}FeF_3$  and both nuclear and magnetic structure of anhydrous HTB-FeF<sub>3</sub>. The influence of frustration on magnetic structure and the role of water for structure stability will have to be investigated.

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